

Extracting flow information from sparse Lagrangian trajectories

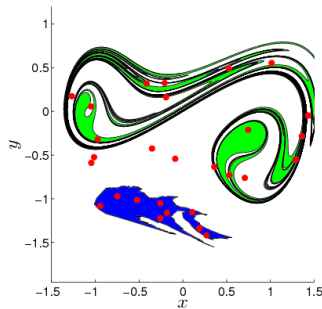
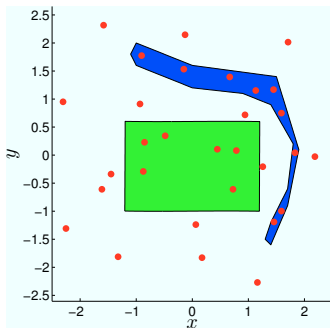
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“Euler Lagrangian meeting,” Wolfgang Pauli Institute, Vienna
7 May 2012

Sparse trajectories and material loops

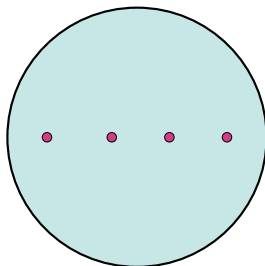


How do we efficiently detect trajectories that 'bunch' together?

[movie 1]

Mathematical background: Punctured disks

Low-dimensional topologists have long studied [transformations of surfaces](#) such as the [punctured disk](#):

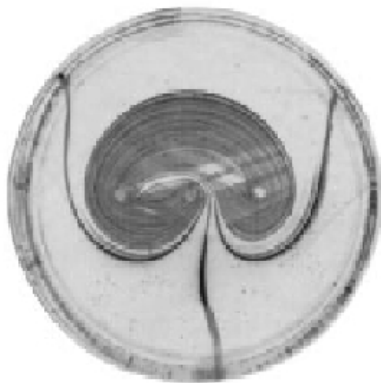


The central object of study is the [homeomorphism](#): a continuous, invertible transformation whose inverse is also continuous.

For instance, this is a model of a two-dimensional vat of viscous fluid with stirring rods.

Punctured disks in experiments

The transformation in this case is given by the solution of a fluid equation over one period of rod motion.

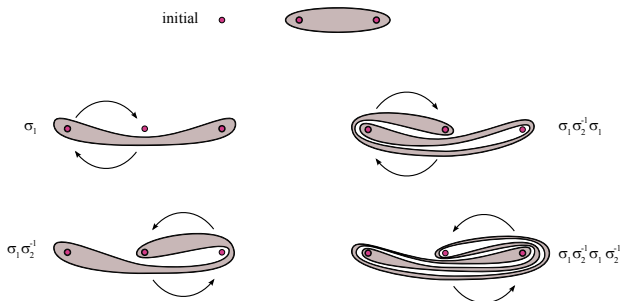


[P. L. Boyland, H. Aref, and M. A. Stremler, *J. Fluid Mech.* **403**, 277 (2000)]

[movie 2] [movie 3]

Growth of curves on a disk

On a disk with 3 punctures (rods), we can also look at the growth of curves:



We use the **braid generator** notation: σ_i means the clockwise interchange of the i th and $(i + 1)$ th rod. (Inverses are counterclockwise.)

The motion above is denoted $\sigma_1 \sigma_2^{-1}$.

Growth of curves on a disk (2)

The rate of growth $h = \log \lambda$ is called the [topological entropy](#).

But how do we find the rate of growth of curves for motions on the disk?

For 3 punctures it's easy: the entropy for $\sigma_1\sigma_2^{-1}$ is $h = \log \varphi^2$, where φ is the [Golden Ratio](#)!

For more punctures, use [Moussafir iterative technique](#) (2006).

[Thiffeault, *Phys. Rev. Lett.* (2005); *Chaos* (2010); Gouillart et al., *Phys. Rev. E* (2006) '[ghost rods](#)']

Iterating a loop

It is well-known that the entropy can be obtained by applying the motion of the punctures to a closed curve (loop) repeatedly, and measuring the growth of the length of the loop (Bowen, 1978).

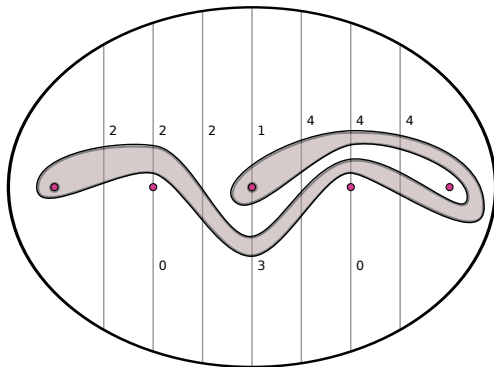
The problem is twofold:

1. Need to keep track of the loop, since its length is growing exponentially;
2. Need a simple way of transforming the loop according to the motion of the punctures.

However, simple closed curves are easy objects to manipulate in 2D. Since they cannot self-intersect, we can describe them **topologically** with very few numbers.

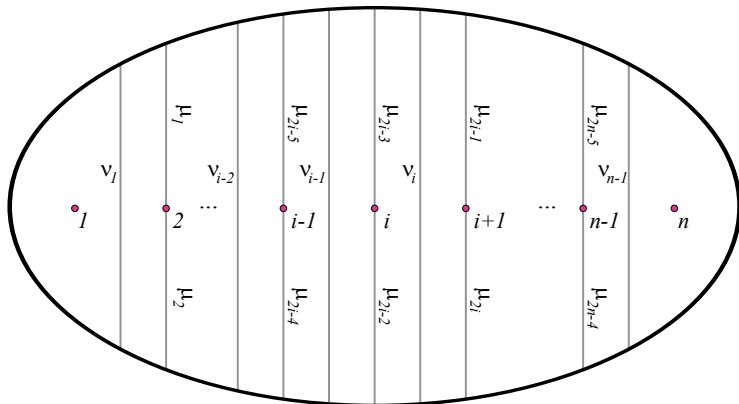
Solution to problem 1: Loop coordinates

What saves us is that a closed loop can be uniquely reconstructed from the number of intersections with a set of curves. For instance, the **Dynnikov coordinates** involve intersections with vertical lines:



Crossing numbers

Label the crossing numbers:



Dynnikov coordinates

Now take the difference of crossing numbers:

$$a_i = \frac{1}{2} (\mu_{2i} - \mu_{2i-1}),$$
$$b_i = \frac{1}{2} (\nu_i - \nu_{i+1})$$

for $i = 1, \dots, n - 2$.

The vector of length $(2n - 4)$,

$$\mathbf{u} = (a_1, \dots, a_{n-2}, b_1, \dots, b_{n-2})$$

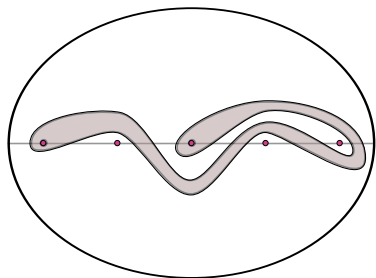
is called the **Dynnikov coordinates** of a loop.

There is a one-to-one correspondence between closed loops and these coordinates: you can't do it with fewer than $2n - 4$ numbers.

Intersection number

A useful formula gives the **minimum intersection number** with the 'horizontal axis':

$$L(\mathbf{u}) = |a_1| + |a_{n-2}| + \sum_{i=1}^{n-3} |a_{i+1} - a_i| + \sum_{i=0}^{n-1} |b_i|,$$

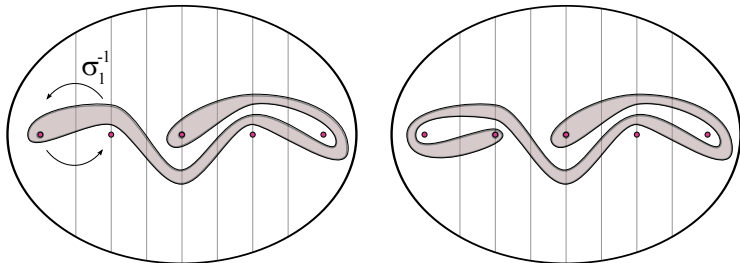


For example, the loop on the left has $L = 12$.

The crossing number grows proportionally to the length.

Solution to problem 2: Action on coordinates

Moving the punctures according to a braid generator changes some crossing numbers:



There is an explicit formula for the change in the coordinates!

Action on loop coordinates

The **update rules** for σ_i acting on a loop with coordinates (\mathbf{a}, \mathbf{b}) can be written

$$a'_{i-1} = a_{i-1} - b_{i-1}^+ - (b_i^+ + c_{i-1})^+,$$

$$b'_{i-1} = b_i + c_{i-1}^-,$$

$$a'_i = a_i - b_i^- - (b_{i-1}^- - c_{i-1})^-,$$

$$b'_i = b_{i-1} - c_{i-1}^-,$$

where

$$f^+ := \max(f, 0), \quad f^- := \min(f, 0).$$

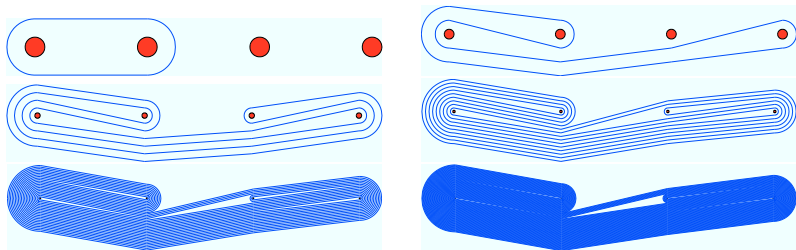
$$c_{i-1} := a_{i-1} - a_i - b_i^+ + b_{i-1}^-.$$

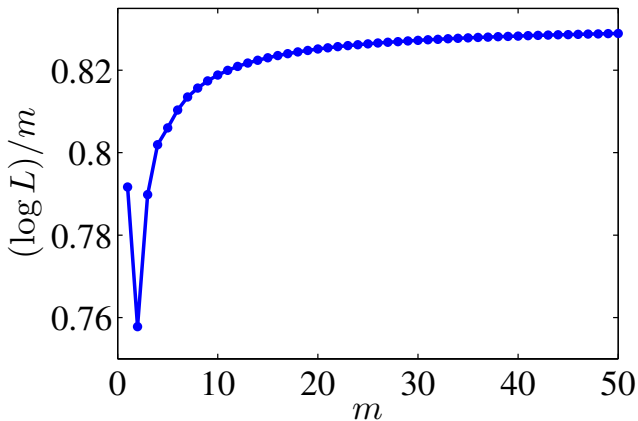
This is called a **piecewise-linear action**.

Easy to code up (see for example Thiffeault (2010)).

Growth of L

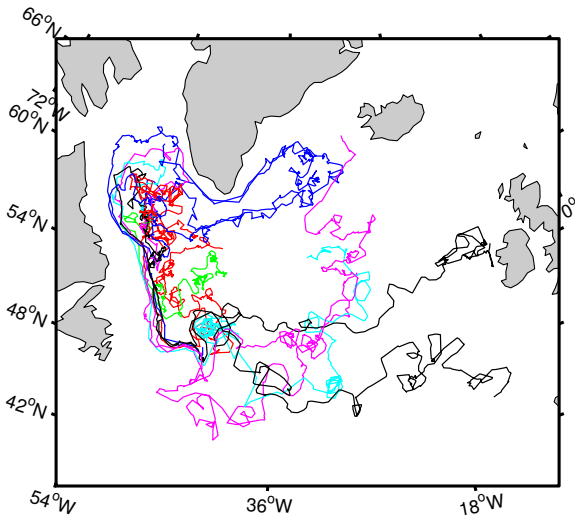
For a specific rod motion, say as given by the braid $\sigma_3^{-1}\sigma_2^{-1}\sigma_3^{-1}\sigma_2\sigma_1$, we can easily see the exponential growth of L and thus measure the entropy:



Growth of L (2)

m is the number of times the braid acted on the initial loop.

Oceanic float trajectories



Oceanic floats: Data analysis

What can we measure?

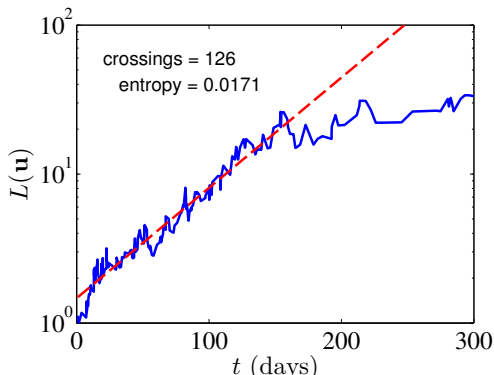
- Single-particle dispersion (not a good use of all data)
- Correlation functions (what do they mean?)
- Lyapunov exponents (some luck needed!)

Another possibility:

Compute the σ_i for the float trajectories (convert to a sequence of symbols), then look at how loops grow. Obtain a **topological entropy** for the motion (similar to Lyapunov exponent).

Oceanic floats: Entropy

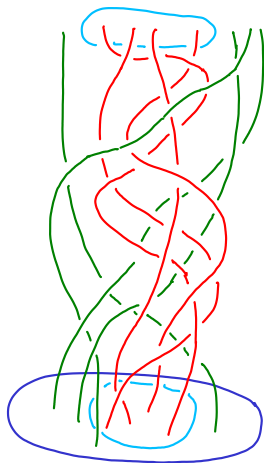
10 floats from Davis' Labrador sea data:



Floats have an entanglement time of about 50 days — timescale for horizontal stirring.

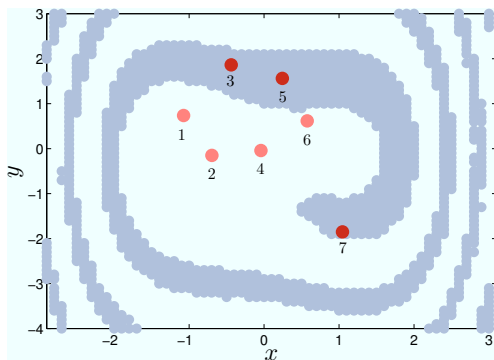
Source: WOCE subsurface float data assembly center (2004)

Lagrangian Coherent Structures



- There is a lot more information in the braid than just entropy;
- For instance: imagine there is an **isolated region** in the flow that does not interact with the rest, bounded by **Lagrangian coherent structures** (LCS);
- Identify LCS and invariant regions from particle trajectory data by searching for curves that grow slowly or not at all.
- For now: regions are not 'leaky.'
- (See the work of Haller et al.)

Sample system: Modified Duffing oscillator

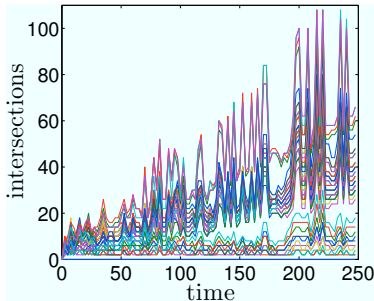
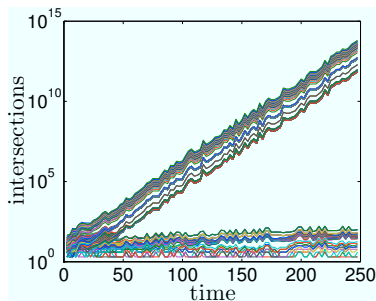


$$\dot{x} = y + \alpha \cos \omega t,$$

$$\dot{y} = x(1 - x^2) + \gamma \cos \omega t - \delta y,$$

+ rotation to further hide two regions. $\alpha = .1$, $\gamma = .14$, $\delta = .08$, $\omega = 1$.

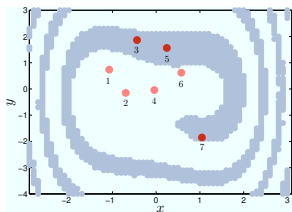
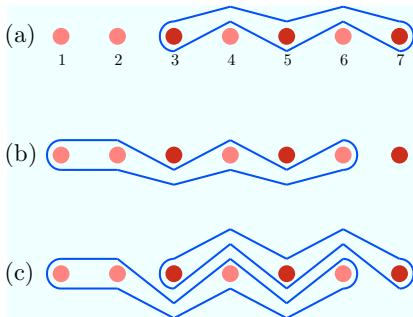
Growth of a vast number of loops



Left: semilog plot; Right: linear plot of slow-growing loops.

Clearly two types of loops!

What do the slowest-growing loops look like?



[(c) appears because the coordinates also encode 'multiloops.']

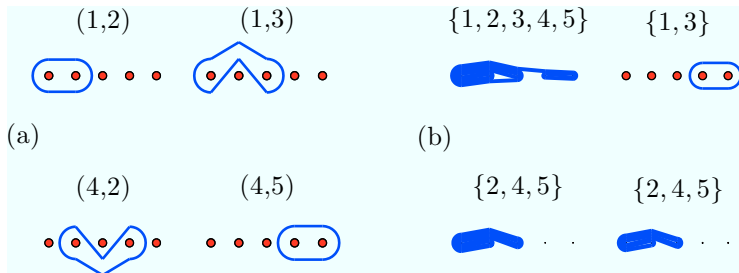
Computational complexity

Here's the bad news:

- There are an infinite number of loops to consider.
- But we don't really expect hyper-convoluted initial loops (nor do we care so much about those).
- Even if we limit ourselves to loops with Dynnikov coordinates between -1 and 1 , this is still 3^{2n-4} loops.
- This is too many... can only treat about 10–11 trajectories using this [direct method](#).

An improved method: Pair-loops

The biggest problem is that we only look at whether a loop grows or not. But there is a lot more information to be found in **how a loop entangles the punctures** as it evolves.



Consider loops that enclose two punctures at once. **More involved analysis, but scales *much* better with n .**

Improvement

Run times in seconds:

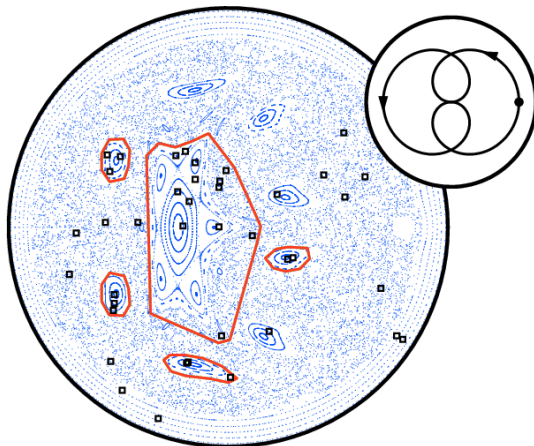
# of trajectories	6	7	8	9	10	11	20
direct method	0.46	0.70	6.0	53	462	3445	N/A
pair-loop method	9.5	11.6	12.3	13	15	20	128

Bottleneck for the pair-loop method is finding the non-growing loops. (Should scale as n^2 for large enough n .)

The downside is that the pair-loop method is much more complicated. But in the end it accomplishes the same thing.

See Allshouse & Thiffeault, *Physica D* **241**, 95–105 (2012).

A physical example: Rod stirring device



[movie 4]

Another benchmark problem: double-gyre

Shadden et al. (2005)

$$\dot{\mathbf{x}} = \pi A \begin{pmatrix} -\sin(\pi f(x, t)) \cos(\pi y) \\ \cos(\pi f(x, t)) \sin(\pi y) \frac{\partial f(x, t)}{\partial x} \end{pmatrix}$$

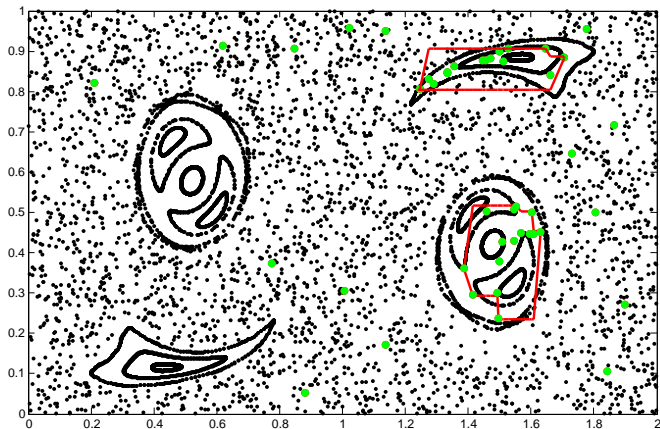
$$f(x, t) = a(t)x^2 + b(t)x$$

$$a(t) = \varepsilon \sin(\omega t)$$

$$b(t) = 1 - 2\varepsilon \sin(\omega t)$$

$$\varepsilon = 0.1, A = 0.1, \omega = \pi/5.$$

Double-gyre coherent structures



[movie 5]

Conclusions

- Having rods undergo ‘braiding’ motion guarantees a minimal amount of entropy ([stretching of material lines](#));
- This idea can also be used on fluid particles to estimate entropy;
- Need a way to compute entropy fast: [loop coordinates](#);
- There is a lot more information in this braid: extract it! ([Lagrangian coherent structures](#));
- Is this useful? We’re still looking for a good data set to try it on! No joy so far. . .
- We’re developing Matlab tools — [braidlab](#).
- Also applicable to [granular media](#) [Puckett \(2012\)](#).
- See [Thiffeault \(2005, 2010\)](#) and [Allshouse & Thiffeault \(2012\)](#).

This work was supported by the Division of Mathematical Sciences of the US National Science Foundation, under grant DMS-0806821.

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